

**Problem 1**

As for any conformal map, it is sufficient to examine the transformation of the boundary of the domain, which in this case is the unit circle:

$$z = \frac{2\zeta}{\zeta - 1} \Rightarrow z(e^{i\theta}) = 2 \frac{\exp(i\theta)}{\exp(i\theta) - 1} \cdot \frac{\exp(-i\theta) - 1}{\exp(-i\theta) - 1} = 2 \frac{1 - \exp(i\theta)}{2 - 2\cos\theta} = \frac{1 - \cos\theta - i\sin\theta}{1 - \cos\theta}$$

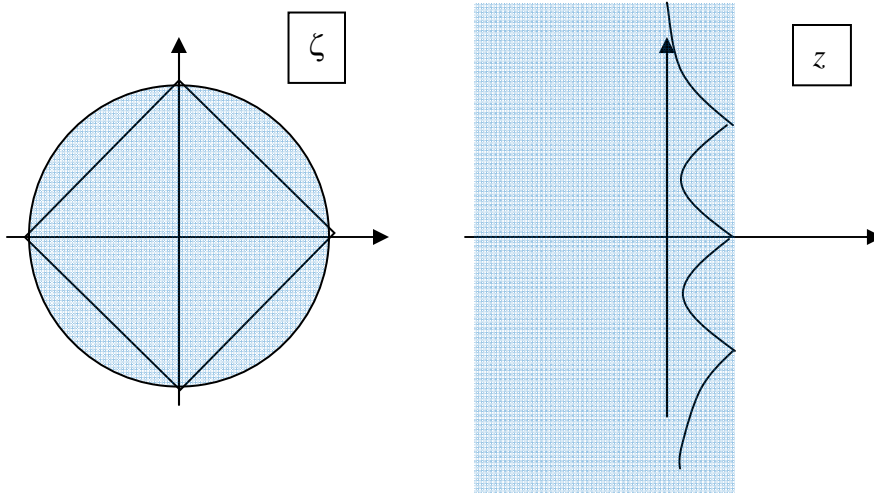
$$\Rightarrow z = 1 - i \frac{\sin\theta}{1 - \cos\theta} \quad \text{Note: we can simplify even further: } z = 1 - i \frac{2\sin\frac{\theta}{2}\cos\frac{\theta}{2}}{2\sin^2\frac{\theta}{2}} = 1 - i \cot\frac{\theta}{2}$$

Another method: 
$$\frac{2\exp(i\theta)}{\exp(i\theta) - 1} = \frac{2\exp(i\theta/2)}{\exp(i\theta/2) - \exp(-i\theta/2)} = \frac{\cos\frac{\theta}{2} + i\sin\frac{\theta}{2}}{i\sin\frac{\theta}{2}} = 1 - i \cot\frac{\theta}{2}$$

So the image of the unit circle is a vertical line  $\text{Re}(z)=1$

Since the center of the disk is transformed into  $z(0)=0$ , the disk is transformed into half-plane to the left of the line  $\text{Re}(z)=1$ .

Transformation of the square:  $z(1) = +\infty$ ,  $z(-1) = -2/(-2) = 1$ ; to get  $z(\pm i)$  set  $\theta$  to  $\pm\pi/2$  in above formula, or calculate directly:  $z(\pm i) = \frac{\pm 2i}{\pm i - 1} = \left( \frac{\pm 2i}{\pm i - 1} \right) \left( \frac{\mp i - 1}{\mp i - 1} \right) = \frac{2 \mp 2i}{2} = 1 \mp i$

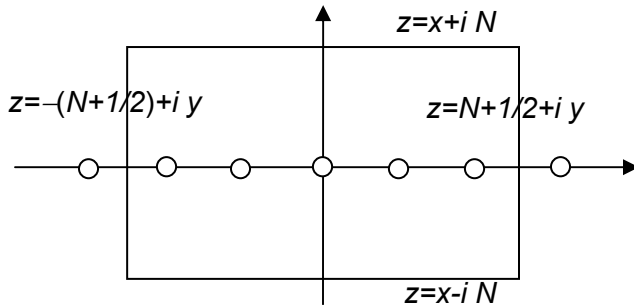


Angles are preserved, so the curved segments in the image of the square meet at right angles

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## Problem 2

When calculating  $\sum_{n=1}^{\infty} \frac{1}{n^k}$  by integrating  $f(z) = \frac{\pi \cos \pi z}{z^k \sin \pi z}$ , the only purpose of the numerator is to cancel the alternating sign coming from the derivative of  $\sin \pi z$  at the pole  $z=n$ . Therefore, we have to get rid of the cosine, and integrate  $f(z) = \frac{\pi}{z^2 \sin \pi z}$  over the same rectangular contour to find  $\sum_{n=1}^{\infty} \frac{(-1)^n}{n^2}$ .



Along horizontal segments  $1/|\sin(\pi z)| \leq 1/|\exp(\pi N) - \exp(-\pi N)|$ , exponentially decaying as  $N = |\text{Im } z| \rightarrow \infty$ , therefore the integral also approaches zero in this limit. Along the vertical segments  $z = \pm(N+1/2) + iy$ , we have

$$\begin{aligned} \frac{1}{|\sin \pi z|} &= \frac{1}{|\sin(\pm(N+\frac{1}{2})\pi + i\pi y)|} = \frac{1}{|\sin(\pm(N+\frac{1}{2})\pi) \cosh \pi y + i \cos(\pm(N+\frac{1}{2})\pi) \sinh \pi y|} = \\ &= \frac{1}{|\sin(\pm(N+\frac{1}{2})\pi) \cosh \pi y|} = \frac{1}{\cosh \pi y} \leq 1 \end{aligned}$$

So the integral is bounded by the length of the contour,  $2N$ , times  $1/|z|^2 \sim 1/N^2$ , yielding zero in the limit  $N \rightarrow \infty$ .

Thus, the sum of residues is zero:

$$\sum_{n=-\infty}^{\infty} \text{Res}(f(z), z_n = n) = \text{Res}(f(z), 0) + 2 \sum_{n=1}^{\infty} \text{Res}(f(z), n) = \text{Res}(f(z), 0) + 2 \sum_{n=1}^{\infty} \frac{(-1)^n}{n^2} \Rightarrow \sum_{n=1}^{\infty} \frac{(-1)^n}{n^2} = -\frac{\text{Res}(f(z), 0)}{2}$$

Now let's find the residue at 0:

$$\begin{aligned} \frac{\pi}{z^2 \sin \pi z} &= \frac{\pi}{z^2 \left( \pi z - \frac{(\pi z)^3}{3!} + O(z^5) \right)} = \frac{1}{z^3} \frac{1}{1 - \frac{(\pi z)^2}{6} + O(z^4)} = \frac{1}{z^3} \left( 1 + \frac{(\pi z)^2}{6} + O(z^4) \right) = \frac{1}{z^3} + \frac{\pi^2}{6} \frac{1}{z} + O(z) \\ \Rightarrow & \boxed{\sum_{n=1}^{\infty} \frac{(-1)^n}{n^2} = -\frac{\pi^2}{12}} \end{aligned}$$

The answer is negative since the first (dominant) term is negative.

### Problem 3

Solve the following Volterra integral equation. You may use without derivation the Laplace transform result

$L[\sin \omega t] = \frac{\omega}{s^2 + \omega^2}$ ; denote  $\gamma = \omega(\omega - 1)$  to simplify algebra. Does solution exist when  $\omega = +1$ ?

$$y(t) = t + \int_0^t y(t - \tau) \sin \omega \tau d\tau$$

$$Y(s) = \frac{1}{s^2} + Y(s) \frac{\omega}{s^2 + \omega^2} \Rightarrow Y(s) \left[ 1 - \frac{\omega}{s^2 + \omega^2} \right] = \frac{1}{s^2} \Rightarrow Y(s) \frac{s^2 + \overbrace{\omega^2 - \omega}^{\gamma = \omega(\omega - 1)}}{s^2 + \omega^2} = \frac{1}{s^2} \Rightarrow Y(s) = \frac{s^2 + \omega^2}{s^2 (s^2 + \gamma)}$$

Since  $Y(s)$  is not that messy, partial fractions can be used:

$$Y(s) = \frac{s^2}{s^2 (s^2 + \gamma)} + \frac{\omega^2}{s^2 (s^2 + \gamma)} = \frac{1}{s^2 + \gamma} + \frac{\omega^2}{\gamma} \left[ \frac{1}{s^2} - \frac{1}{s^2 + \gamma} \right] = \frac{\omega^2}{\gamma} \frac{1}{s^2} + \underbrace{\left( 1 - \frac{\omega^2}{\gamma} \right)}_{-\omega \gamma} \frac{1}{s^2 + \gamma} = \boxed{\frac{\omega^2}{\gamma} \frac{1}{s^2} - \frac{\omega}{\gamma} \frac{1}{s^2 + \gamma}}$$

Since  $L[\sin \omega t] = \frac{\omega}{s^2 + \omega^2}$ , we infer that  $L^{-1} \left[ \frac{1}{s^2 + \gamma} \right] = \frac{1}{\sqrt{\gamma}} \sin(\sqrt{\gamma} t)$  therefore:

$$\Rightarrow y(t) = \frac{\omega^2}{\gamma} t - \frac{\omega}{\gamma^{3/2}} \sin \sqrt{\gamma} t = \boxed{\frac{\omega}{\omega - 1} t - \frac{1}{\sqrt{\omega(\omega - 1)^3}} \sin \sqrt{\omega(\omega - 1)} t}$$

Singularity  $\omega=1$  is removable:  $\gamma=0 \Rightarrow Y(s) = \frac{s^2 + 1}{s^2 (s^2 + 0)} = \frac{1}{s^2} + \frac{1}{s^4} \Rightarrow \boxed{Y(s) = t + \frac{t^3}{6}}$

(Note: for  $0 < \omega < 1$  the above results are still correct, yielding  $y(t) = \frac{\omega}{\omega - 1} t + \frac{1}{\sqrt{\omega(1 - \omega)^3}} \sinh \sqrt{\omega(1 - \omega)} t$ )

#### Problem 4

Use the Fourier transform to write down the solution for the following modification of the heat (or diffusion) equation

$$\begin{cases} \frac{\partial u}{\partial t} = \frac{\partial^2 u}{\partial x^2} - \omega \frac{\partial u}{\partial x}; & -\infty < x < +\infty, \quad t > 0, \\ u(x, t = 0) = f(x) \end{cases} \Rightarrow \begin{cases} \frac{\partial U(k, t)}{\partial t} = (-k^2 - ik\omega) U(k, t) \\ U(k, 0) = F(k) \end{cases}$$

$$\Rightarrow U(k, t) = F(k) \underbrace{\exp(-(k^2 + ik\omega)t)}_{G(k)} \Rightarrow u(x, t) = \int_{-\infty}^{+\infty} f(x - \xi) g(\xi, t) d\xi$$

where the Green's function is

$$\begin{aligned} g(x, t) &= \frac{1}{2\pi} \int_{-\infty}^{+\infty} \underbrace{\exp(-(k^2 + ik\omega)t)}_{G(k)} \exp(ikx) dk = \frac{1}{2\pi} \int_{-\infty}^{+\infty} \exp(-k^2 t + ik(x - \omega t)) dk \\ &= \frac{1}{2\pi} \int_{-\infty}^{+\infty} \exp\left(-t \left[ k^2 + ik \frac{x - \omega t}{t} \right]\right) dk = \frac{1}{2\pi} \int_{-\infty}^{+\infty} \exp\left(-t \left[ \left( k + i \frac{x - \omega t}{2t} \right)^2 + \left( \frac{x - \omega t}{2t} \right)^2 \right]\right) dk \\ &= \frac{\exp\left(-\frac{(x - \omega t)^2}{4t}\right)}{2\pi} \int_{-\infty}^{+\infty} \exp\left(-t \left( \underbrace{k + i \frac{x - \omega t}{2t}}_s \right)^2\right) dk = \frac{\exp\left(-\frac{(x - \omega t)^2}{4t}\right)}{2\pi} \underbrace{\int_{-\infty + i \frac{x - \omega t}{2t}}^{+\infty + i \frac{x - \omega t}{2t}} \exp(-ts^2) ds}_{\sqrt{\frac{\pi}{t}}} \\ &\Rightarrow \boxed{g(x, t) = \frac{\exp\left(-\frac{(x - \omega t)^2}{4t}\right)}{\sqrt{4\pi t}}} \end{aligned}$$

Note: same Green's function as in the usual diffusion equation, but with the peak moving with velocity  $\omega$

As time evolves, the spread of the Green's function increases, and its center moves to the right

The extra term describes the flow of the medium to the right (transport / advection), with velocity  $\omega$ .